

# Benchmark calculations for electron scattering on atomic hydrogen and helium

I. Bray and D. V. Fursa

ARC Centre for Antimatter-Matter Studies, Curtin University, Perth, Western Australia

IAEA, Vienna, May, 2011



# Outline

- 1 Introduction
  - Motivation
  - Distant history
  - Recent history
- 2 Nonrelativistic CCC theory
  - atomic target structure
  - electron-atom scattering
- 3 Electron-hydrogen scattering
- 4 Electron-helium scattering

# Introduction

## Motivation

- The primary motivation is to provide accurate atomic collision data for science and industry:
  - Astrophysics
  - Fusion research
  - Lighting industry
  - Medical and materials applications
- Provide a rigorous foundation for collision theory with long-ranged (Coulomb) potentials.

# Introduction

## Motivation

- The primary motivation is to provide accurate atomic collision data for science and industry:
  - Astrophysics
  - Fusion research
  - Lighting industry
  - Medical and materials applications
- Provide a rigorous foundation for collision theory with long-ranged (Coulomb) potentials.

# Introduction

## Distant history

- Prior to the 1990s theory and experiment generally did not agree for:
  - electron-hydrogen excitation or ionisation,
  - electron-helium excitation or ionisation,
  - single or double photoionisation of helium.
- Consequently, we have been developing the convergent close-coupling (CCC) theory for electron/positron/photon/(anti)proton collisions with atoms/ions/molecules that is applicable at all energies for the major excitation and ionisation processes.

# Introduction

## Distant history

- Prior to the 1990s theory and experiment generally did not agree for:
  - electron-hydrogen excitation or ionisation,
  - electron-helium excitation or ionisation,
  - single or double photoionisation of helium.
- Consequently, we have been developing the convergent close-coupling (CCC) theory for electron/positron/photon/(anti)proton collisions with atoms/ions/molecules that is applicable at all energies for the major excitation and ionisation processes.

# Introduction

## Recent history

- Prior to 2008, no satisfactory mathematical formulation in the case of long-range (Coulomb) potentials for positive-energy scattering in
  - Two-body problems
  - Three-body problems
- Consequently, have developed a surface integral approach to scattering theory that is valid for short- and long-range potentials, see Kadyrov et al. PRL, **101**, 230405 (2008) and Annals of Physics **324** 1516 (2009).



# Introduction

## Recent history

- Prior to 2008, no satisfactory mathematical formulation in the case of long-range (Coulomb) potentials for positive-energy scattering in
  - Two-body problems
  - Three-body problems
- Consequently, have developed a surface integral approach to scattering theory that is valid for short- and long-range potentials, see Kadyrov et al. PRL, **101**, 230405 (2008) and Annals of Physics **324** 1516 (2009).

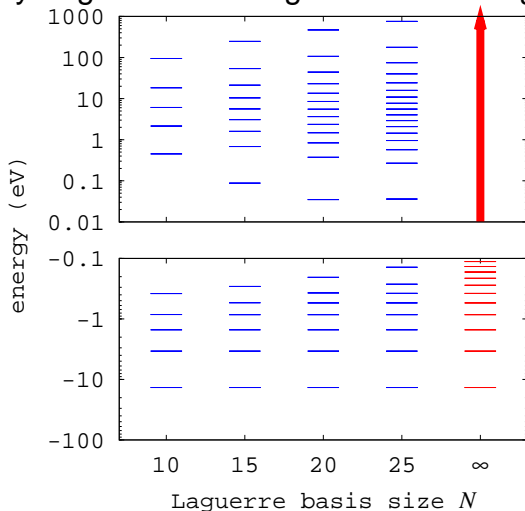
# Nonrelativistic CCC theory

## target structure and scattering

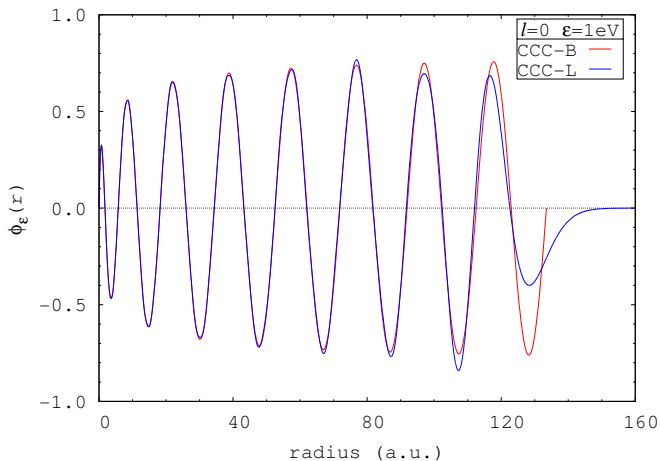
- Structure: using the Laguerre basis  $\xi_{nl}^{(\lambda)}(r)$  write:
  - “one-electron” (H, Li, . . . , Cs) states:  
$$\phi_{nl}^{(\lambda)}(r) = \sum_{n'} C_{nl}^{n'} \xi_{n'l}^{(\lambda)}(r)$$
  - “two-electron” (He, Be, . . . , Hg) states:  
$$\phi_{nls}^{(\lambda)}(r_1, r_2) = \sum_{n', n''} C_{nls}^{n' n''} \xi_{n'l'}^{(\lambda)}(r_1) \xi_{n''l''}^{(\lambda)}(r_2).$$
  - Coefficients  $C$  are obtained by diagonalising the target (FCHF) Hamiltonian

$$\langle \phi_f^{(\lambda)} | H_T | \phi_i^{(\lambda)} \rangle = \varepsilon_f^{(\lambda)} \delta_{fi}. \quad (1)$$

- Hydrogen  $\ell = 0$  energies for  $\lambda = 1$  Laguerre bases



- $\phi_\varepsilon^{(\lambda)}(r)$  for  $N = 70, \lambda = 2$  and  $\phi_\varepsilon^{(R_0)}(r)$  for  $R_0 = 134$ .



- Scattering:

- Electron-atom wavefunction is expanded as

$$|\Psi_i^{(+)}\rangle \approx \mathcal{A} \sum_{n=1}^N |\phi_n^{(\lambda)}\rangle \langle \phi_n^{(\lambda)} | \psi_i^{(+)}\rangle \equiv \mathcal{A} I_N |\psi_i^{(+)}\rangle. \quad (2)$$

- Solve for  $T_{fi} \equiv \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \Psi_i^{(+)}\rangle$  at  $E = \varepsilon_i^{(\lambda)} + \epsilon_{k_f}$ ,

$$\begin{aligned} \langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle &= \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \\ &+ \sum_{n=1}^N \int d^3k \frac{\langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_n^{(\lambda)} \mathbf{k} \rangle \langle \mathbf{k} \phi_n^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle}{E + i0 - \varepsilon_n^{(\lambda)} - \epsilon_k}. \end{aligned} \quad (3)$$

- Have step-function behaviour:

$$\langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \approx 0 \text{ for } \varepsilon_f^{(\lambda)} > \epsilon_{k_f}.$$

- Scattering:

- Electron-atom wavefunction is expanded as

$$|\Psi_i^{(+)}\rangle \approx \mathcal{A} \sum_{n=1}^N |\phi_n^{(\lambda)}\rangle \langle \phi_n^{(\lambda)} | \psi_i^{(+)} \rangle \equiv \mathcal{A} I_N |\psi_i^{(+)}\rangle. \quad (2)$$

- Solve for  $T_{fi} \equiv \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \Psi_i^{(+)} \rangle$  at  $E = \varepsilon_i^{(\lambda)} + \epsilon_{k_f}$ ,

$$\begin{aligned} \langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle &= \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \\ &+ \sum_{n=1}^N \iint d^3k \frac{\langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_n^{(\lambda)} \mathbf{k} \rangle \langle \mathbf{k} \phi_n^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle}{E + i0 - \varepsilon_n^{(\lambda)} - \epsilon_k}. \end{aligned} \quad (3)$$

- Have step-function behaviour:

$$\langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \approx 0 \text{ for } \varepsilon_f^{(\lambda)} > \epsilon_{k_f}.$$

- Scattering:

- Electron-atom wavefunction is expanded as

$$|\Psi_i^{(+)}\rangle \approx \mathcal{A} \sum_{n=1}^N |\phi_n^{(\lambda)}\rangle \langle \phi_n^{(\lambda)} | \psi_i^{(+)} \rangle \equiv \mathcal{A} I_N |\psi_i^{(+)}\rangle. \quad (2)$$

- Solve for  $T_{fi} \equiv \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \Psi_i^{(+)} \rangle$  at  $E = \varepsilon_i^{(\lambda)} + \epsilon_{k_f}$ ,

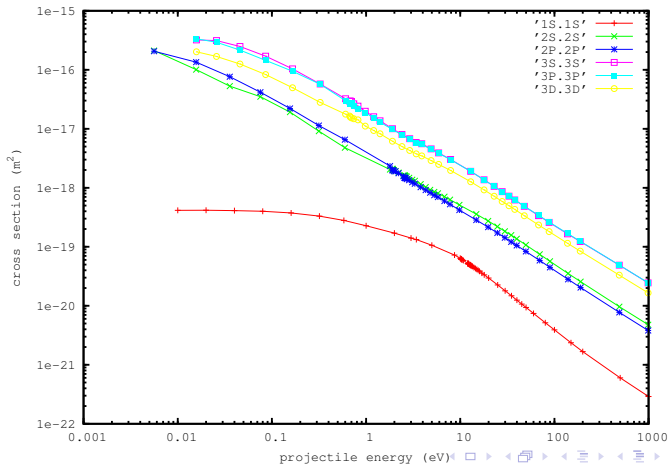
$$\begin{aligned} \langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle &= \langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \\ &+ \sum_{n=1}^N \iint d^3k \frac{\langle \mathbf{k}_f \phi_f^{(\lambda)} | V | \phi_n^{(\lambda)} \mathbf{k} \rangle \langle \mathbf{k} \phi_n^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle}{E + i0 - \varepsilon_n^{(\lambda)} - \epsilon_k}. \end{aligned} \quad (3)$$

- Have step-function behaviour:

$$\langle \mathbf{k}_f \phi_f^{(\lambda)} | T | \phi_i^{(\lambda)} \mathbf{k}_i \rangle \approx 0 \text{ for } \varepsilon_f^{(\lambda)} > \epsilon_{k_f}.$$

# Electron-hydrogen scattering

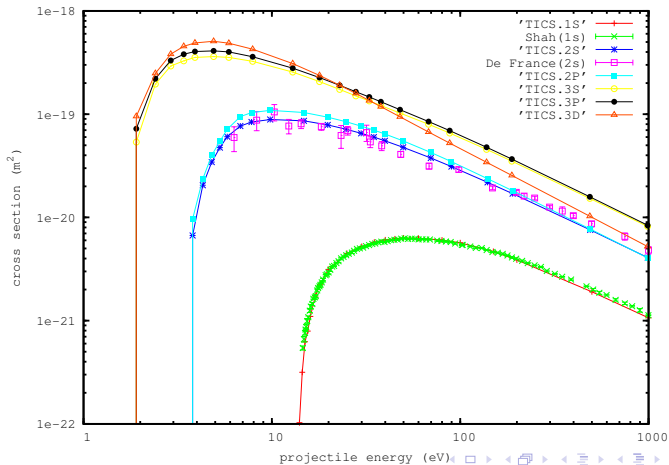
- elastic scattering for  $n \leq 3$  states





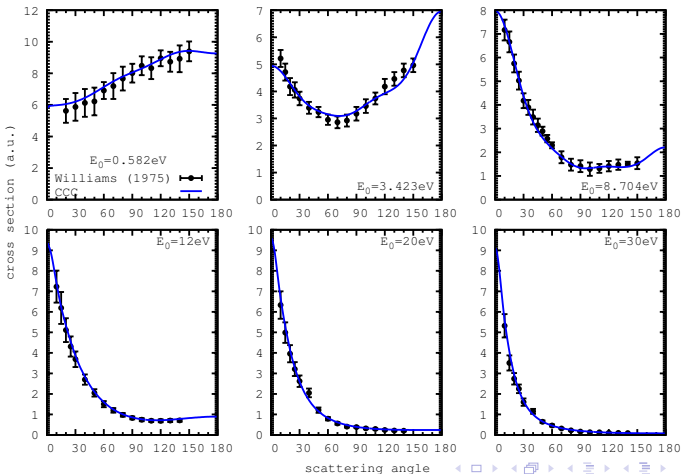
# Electron-hydrogen scattering

- total ionization for  $n \leq 3$  states



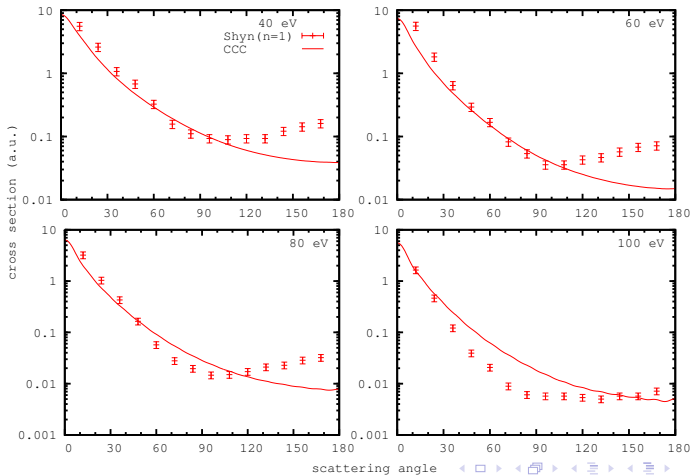
# Electron-hydrogen scattering

- elastic differential cross sections



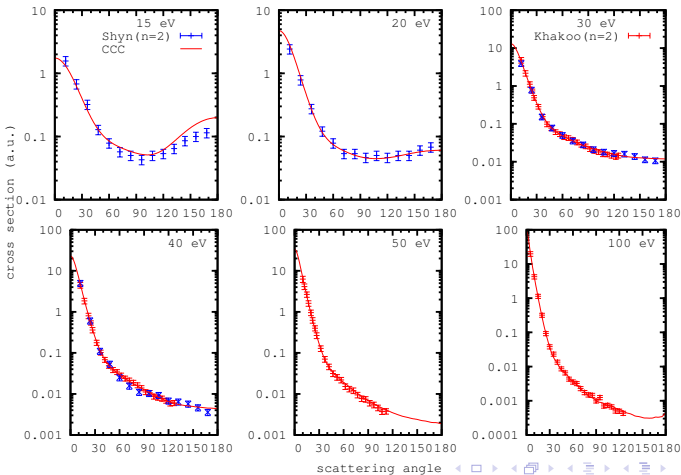
# Electron-hydrogen scattering

- elastic differential cross sections



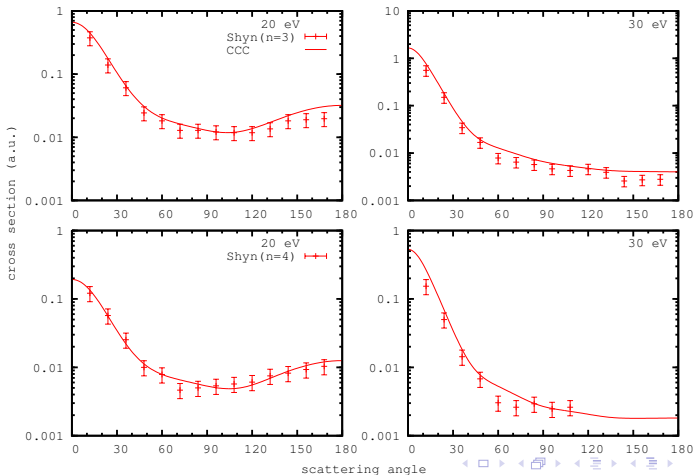
# Electron-hydrogen scattering

## • $n = 2$ excitation differential cross sections



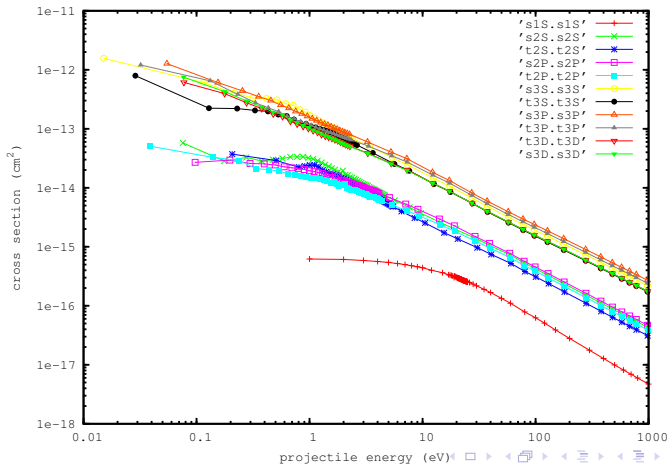
# Electron-hydrogen scattering

- $n = 3, 4$  excitation differential cross sections



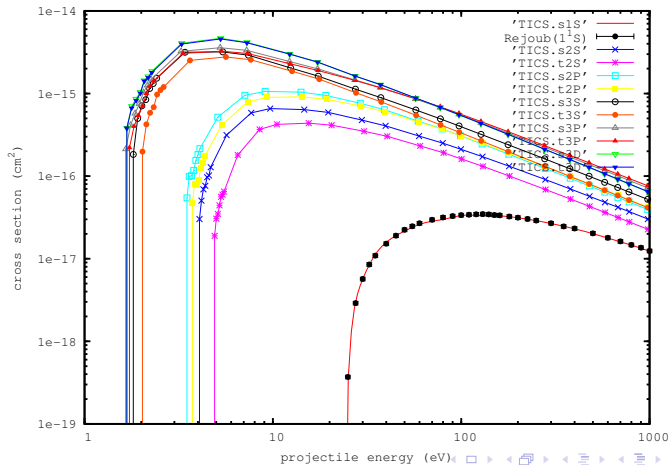
# Electron-helium scattering

- elastic scattering for  $n \leq 3$  states



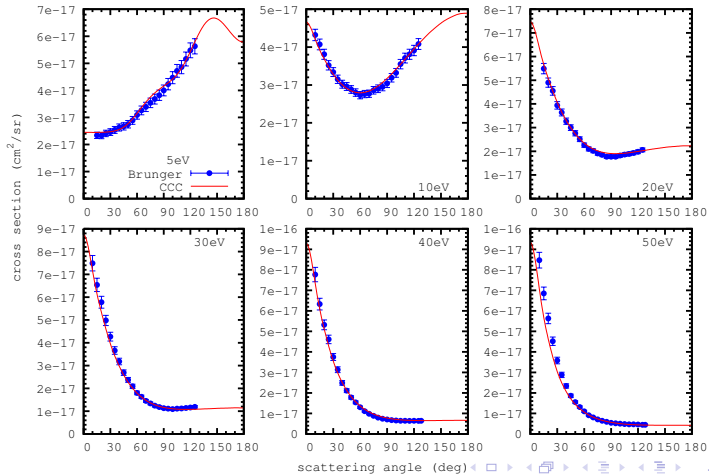
# Electron-helium scattering

- total ionization for  $n \leq 3$  states



# Electron-helium scattering

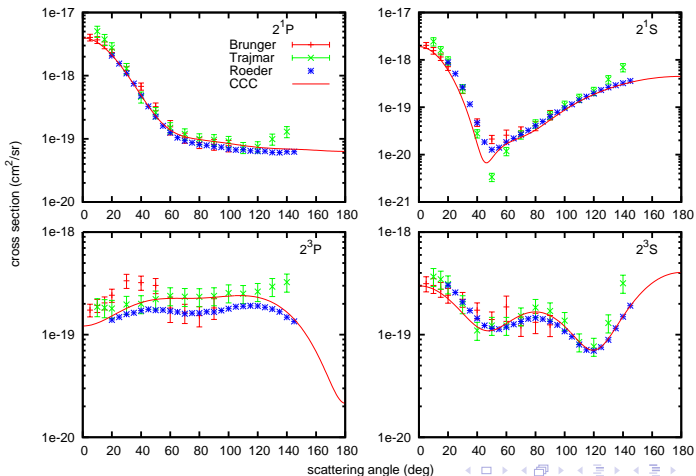
## • elastic differential cross sections





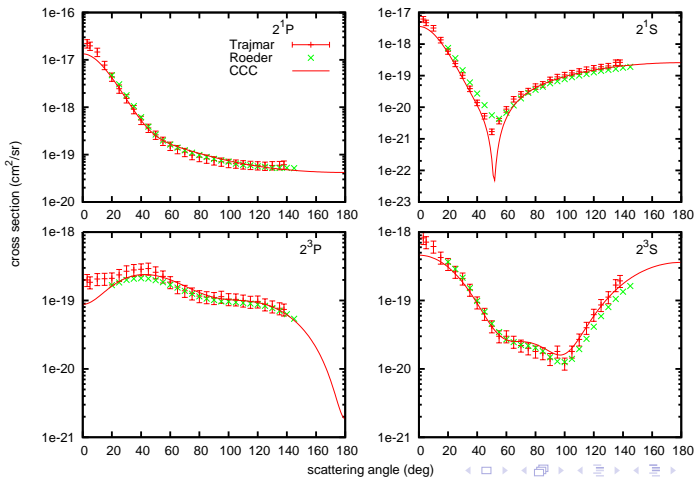
# Electron-helium scattering

## • 30 eV differential cross sections



# Electron-helium scattering

## 40 eV differential cross sections



# Concluding remarks

- There are (almost) no substantial discrepancies between CCC and experiment for:
  - electron-impact excitation or ionisation of quasi one- and two-electron targets
  - photon-impact single and double ionisation of He
- Presently we are extending CCC to
  - positron collisions with quasi one- and two-electron targets
  - (anti)proton collisions with H and He, including ionisation
  - more complicated targets, such as inert gases and molecules.

## Concluding remarks

- There are (almost) no substantial discrepancies between CCC and experiment for:
  - electron-impact excitation or ionisation of quasi one- and two-electron targets
  - photon-impact single and double ionisation of He
- Presently we are extending CCC to
  - positron collisions with quasi one- and two-electron targets
  - (anti)proton collisions with H and He, including ionisation
  - more complicated targets, such as inert gases and molecules.