

Generalized collisional radiative model for light elements: Boron

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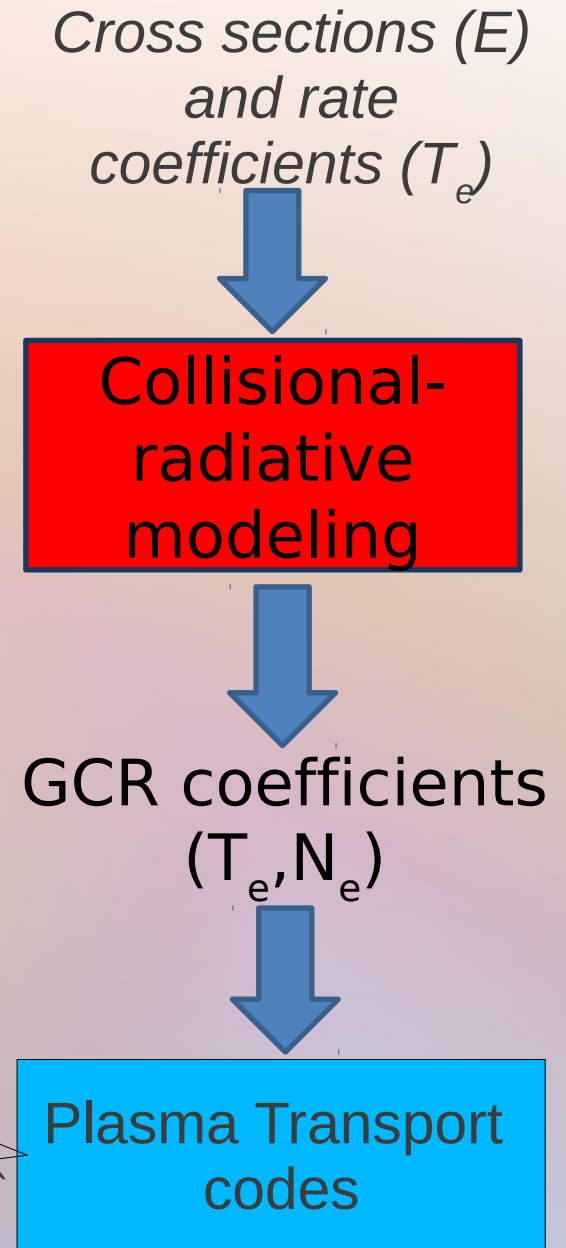
This work was supported by grants from the US Department of Energy. The computational work was carried out on the NERSC and NICS supercomputers.

Outline

- Brief description of
 - Generalized Collisional Radiative theory (GCR)
 - The need for generalized atomic data in fusion
- Examples of GCR data
- New GCR data for boron
 - Description of the fundamental collision data
 - Overview of generating the GCR data
 - Examples of new GCR data
- What remains for GCR work on light species

The GCR approach for fusion applications

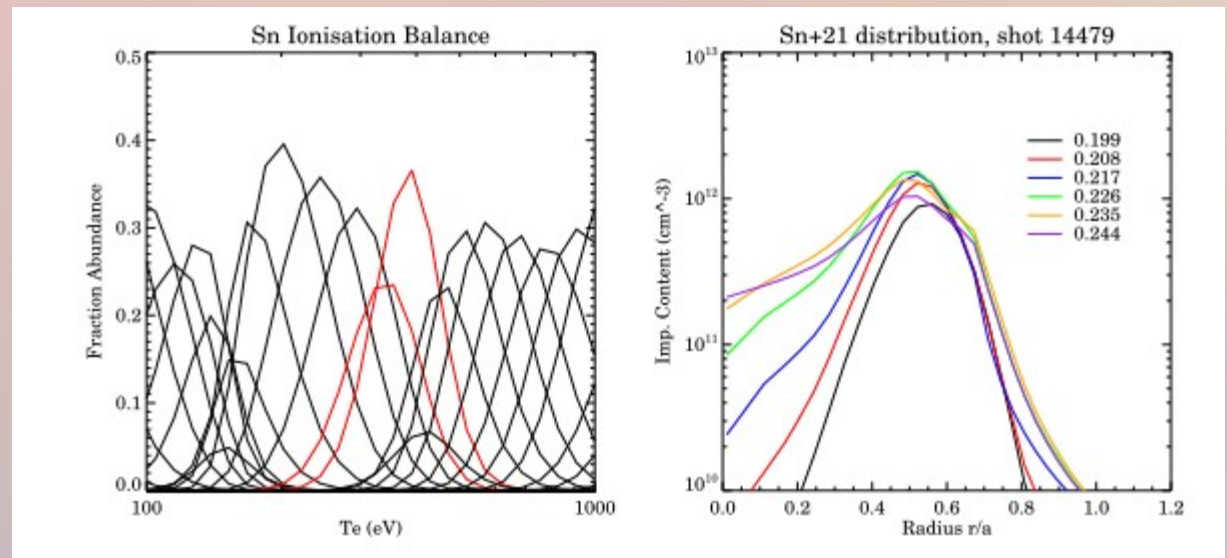
- The fundamental atomic data is processed through a collisional-radiative model to produce data that can be easily used in plasma impurity transport codes. The data is used to track
 - the fractional abundance of the element as it transports in the plasma
 - The radiative power loss (electron cooling)



The importance of GCR data in impurity transport modelling

- Impurity transport codes for fusion (SOLPS, TRANSP, SANCO etc) model the ion stage distribution of impurity species throughout the plasma.
- Both the **ground** and **metastable** populations must be tracked.
- The role of the excited states in the coefficients that connect the ion stages was also found to be important.

$$\frac{dN_z(\rho, t)}{dt} = -\nabla\Gamma_z(\rho, t) - S_{z \rightarrow z+1}N_z(\rho, t) + S_{z-1 \rightarrow z}N_{z-1}(\rho, t) - \alpha_{z \rightarrow z-1}N_z(\rho, t) + \alpha_{z+1 \rightarrow z}N_{z+1}(\rho, t)$$



Tin impurity transport for the MAST experiment, taken from PhD thesis of Foster (2008)

The GCR coefficients

- Ionization

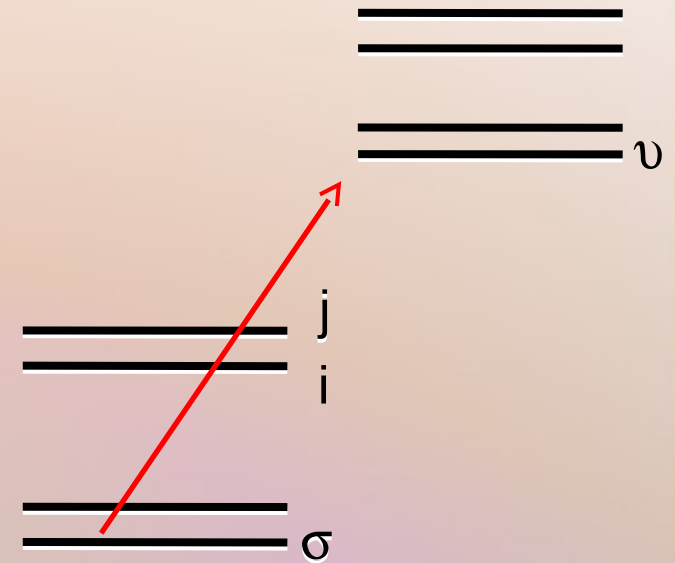
$$S_{CD,\sigma \rightarrow \nu} = (\mathcal{I}_{\nu\sigma} - \sum_{j=1}^O \mathcal{I}_{\nu j} \sum_{i=1}^O \mathcal{C}_{ji}^{-1} \mathcal{C}_{i\sigma})$$

- Recombination

$$\alpha_{CD,\nu' \rightarrow \rho} = (\mathcal{R}_{\rho\nu'} + \sum_{j=1}^O \mathcal{C}_{\rho j} \sum_{i=1}^O \mathcal{C}_{ji}^{-1} \mathcal{R}_{i\nu'})$$

- Photon emissivity

$$P_{LT,\sigma} = \sum_{k,j} \Delta E_{kj} A_{j \rightarrow k} \mathcal{F}_{j\sigma}^{(exc)}$$

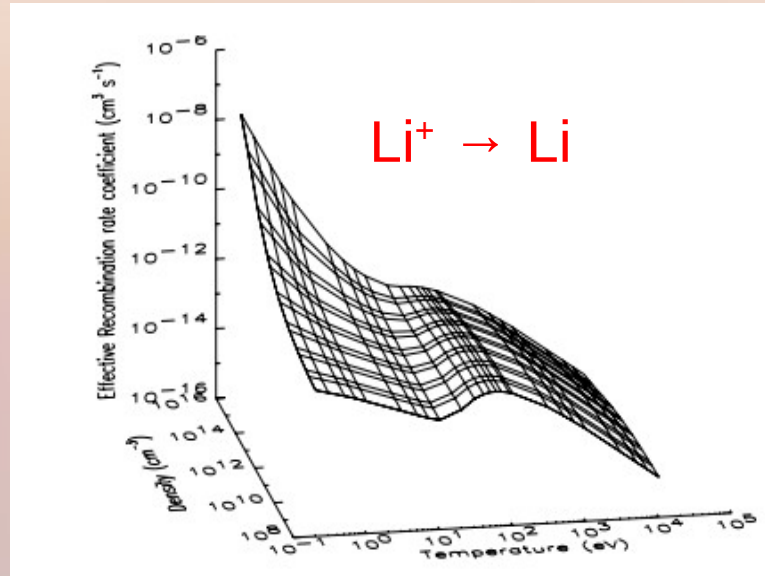


Ionizations per photon for impurity influx diagnostics

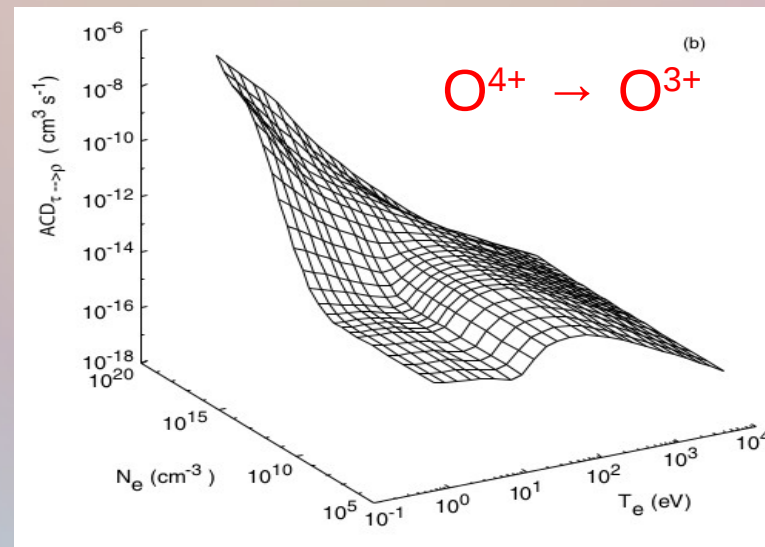
$$SXB_{i-j}^z = \frac{S^{z \rightarrow z+1}(Ne, Te)}{A_{i \rightarrow j} \frac{N_i}{N_z}(Ne, Te)}$$

Examples: GCR recombination rate coefficient

- At low densities DR dominates.
- At moderate densities, collisions reduce the DR (due to some of the excited states being collisionally ionized before it can radiate to the ground/metastable levels).
- At the highest densities, 3-body recombination takes over.



$1s^2 2s (^2S) \rightarrow 1s^2 (^1S)$
 Loch et al.,
 ADNDT,
92 813 (2006)



$2s2p^2 (^4P) \rightarrow 2s^2 (^1S)$
 Summers et al.
 PPCF,**48** 263
 (2006)

GCR ionization

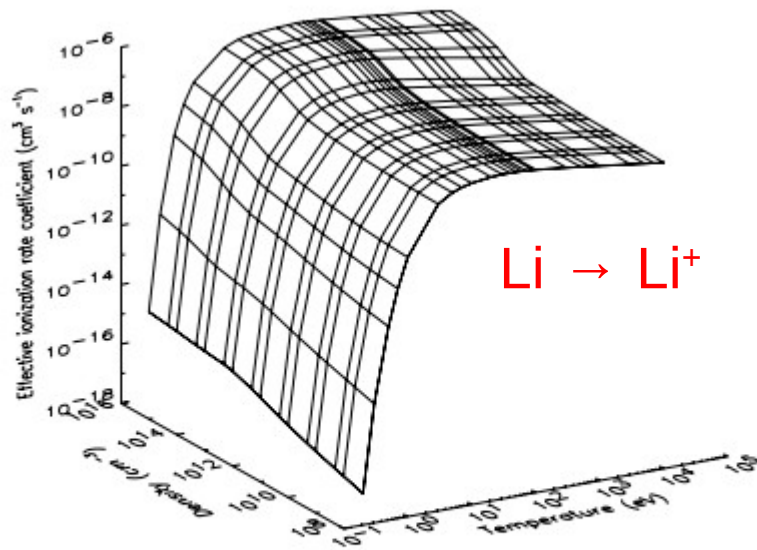


Fig. 8. Effective ionization rate coefficient for the ionization process $e + \text{Li} (1s^2 2s^2 S) \rightarrow \text{Li}^+ (1s^2 2S) + 2e$ as a function of electron temperature and density. Note that the density dependence comes in through the role of ionization from excited states.

Loch et al. ADNDT **92** 818 (2006)

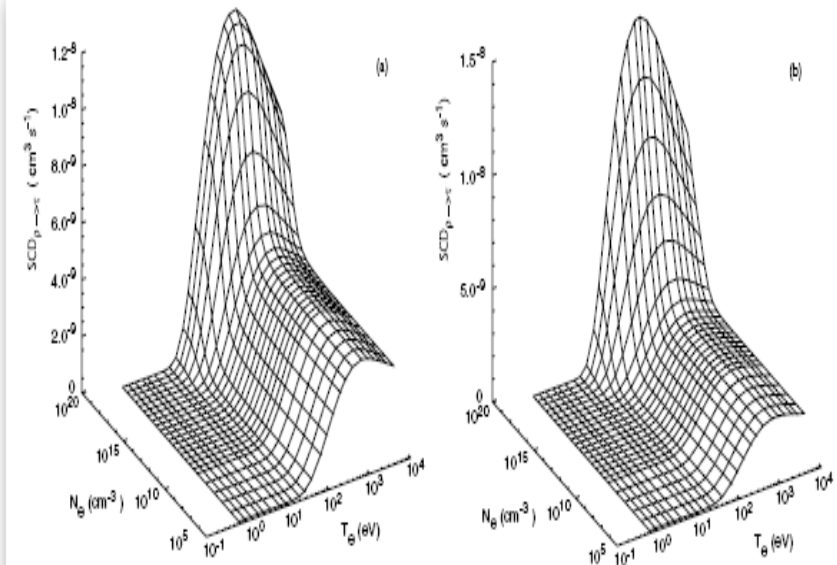
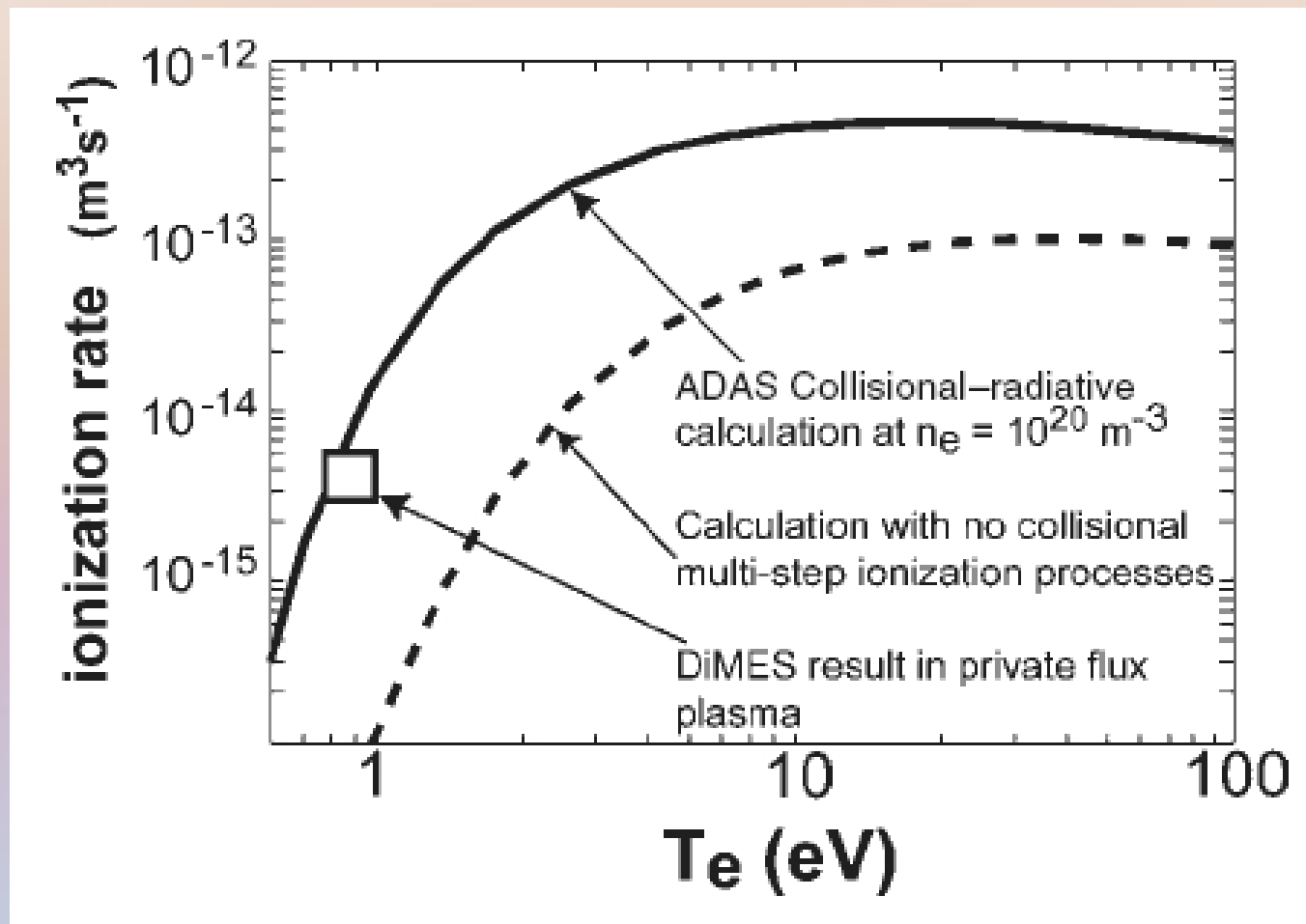


Figure 9. Generalized collisional-radiative ionization coefficients for $\text{O}^{+3} + e \rightarrow \text{O}^{+4} + e + e$ (a) $\text{SCD}(2s^2 2p^2 P \rightarrow 2s^2 1S)$ (b) $\text{SCD}(2s^2 2p^2 P \rightarrow 2s 2p^3 P)$.

Summers et al. PPCF, **48** 263 (2006)

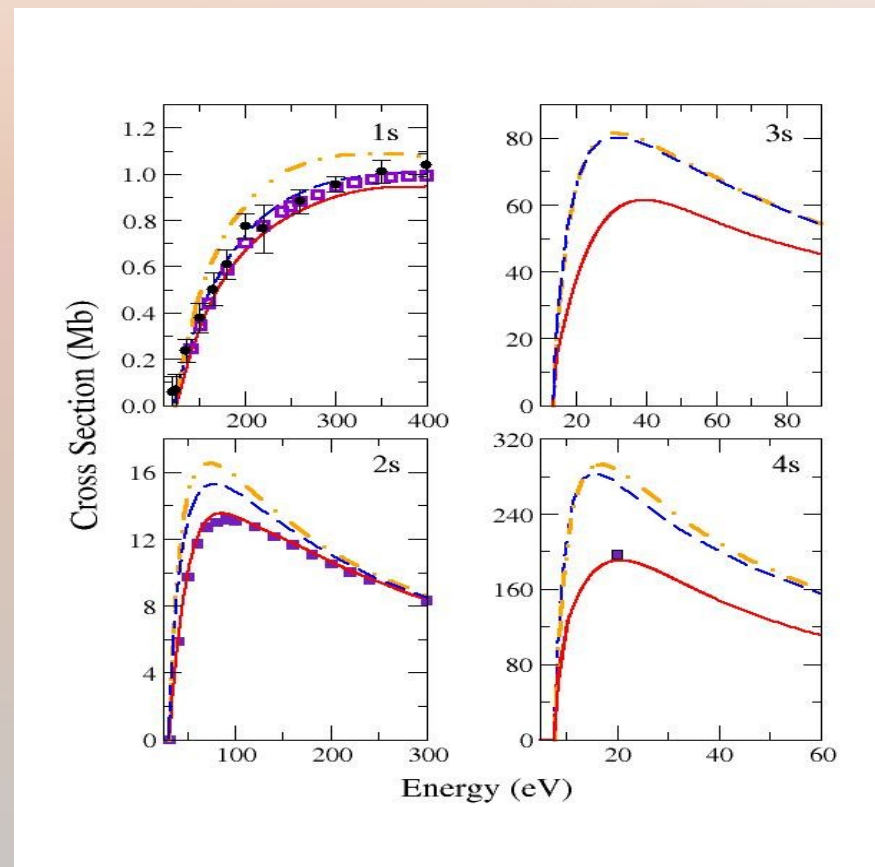
Measurements of Li GCR ionization of the DIII-D tokamak



Allain et al., Nucl. Fusion,
44 655 (2004)

The problem of ionization from excited states

- So one needs data for ionization from the excited levels. However,
- Perturbative methods overestimate the ionization cross section for near neutral systems. *This gets worse for excited states.*
- Calculations using non-perturbative methods (TDCC, RMPS, CCC) become increasingly difficult for higher n-shells.
- There is a need to calculate data up to quite high n-shells.



Griffin et al., *J. Phys. B*, **38** L199 (2005)

Excited states ionization of neutral Boron

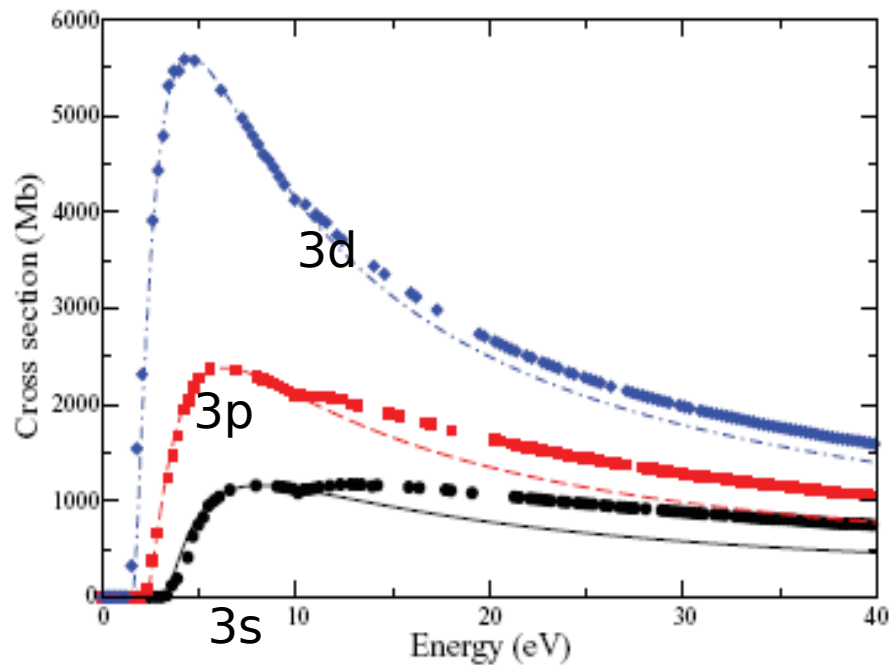


FIG. 1. (Color online) Total electron-impact-ionization cross sections for the $3l$ excited states of B. Circles, raw RMPS for $1s^2 2s^2 3s$; squares, raw RMPS for $1s^2 2s^2 3p$; diamonds, raw RMPS for $1s^2 2s^2 3d$. Solid line, fit to low-energy raw RMPS data for $1s^2 2s^2 3s$; dashed line, fit to low-energy raw RMPS data for $1s^2 2s^2 3p$; dot-dashed line, fit to low-energy raw RMPS data for $1s^2 2s^2 3d$ (1 Mb = 10^{-18} cm²).

Lee et al., *Phys. Rev. A* **82** 042721 (2010)

- Consider the ionization cross sections (RMPS) for the $n=3$ shell in neutral B.
 - Excitation-autoionization starts to contribute above about 10 eV and becomes smaller for the higher n -shells.
 - By fitting the direct ionization part we can see if there is an n -scaling in the cross sections.
 - If it was a purely classical calculation the scaling would go as n^4 .
- We repeated the same study for B^+ , and B^{2+} .

n-scaling data for B, B[±] and B²⁺

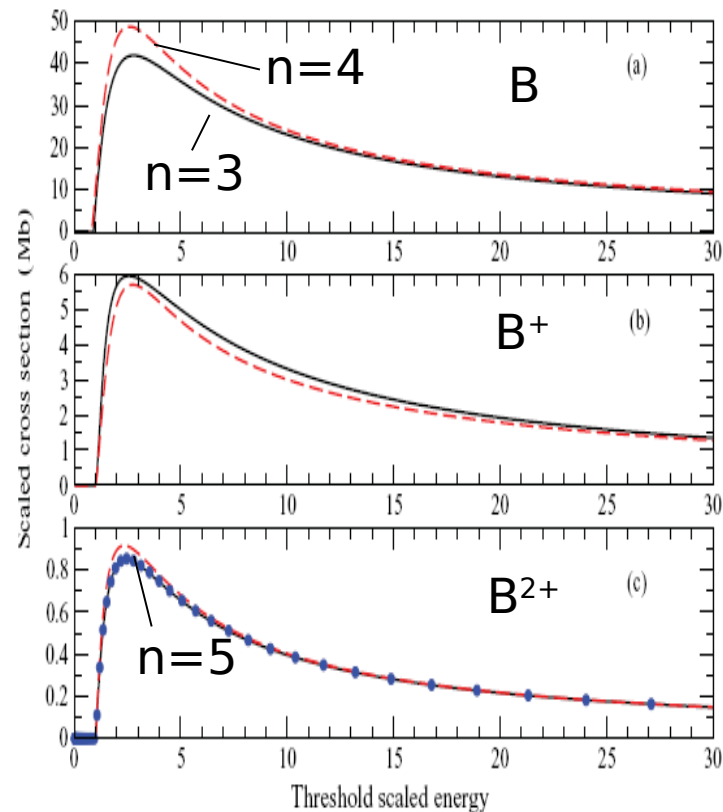


FIG. 5. (Color online) n -scaled electron-impact-ionization cross sections vs threshold scaled energy, that is, cross section divided by n^4 for the n -bundled excited states of (a) B, (b) B⁺, and (c) B²⁺. In all plots the solid line shows the $n = 3$ RMPS data, the dashed line shows the $n = 4$ RMPS data and in panel (c) the solid circles show the $n = 5$ RMPS data (1 Mb = 10⁻¹⁸ cm²).

- For each of the ions a scaling very close to n^4 was found.
- Evaluate your non-perturbative calculation until scales as n^4 , then extrapolate to higher n .
- Or you can fit semi-empirical data (e.g. ECIP) to the RMPS results and used the same scaling factor to scale to even higher n shells.
- Note that the bundled- n , or the spin-resolved data can be extrapolated.

Lee et al., Phys. Rev. A **82** 042721 (2010)

GCR data

- GCR data for the light elements had already been generated within ADAS, using a range of atomic data.
- We have been going through each element and updating the atomic data (if needed), then generating new GCR coefficients:
 - H & He : *Loch et al., Plasma Phys. Control. Fusion, **51** 105006 (2009)*
 - Li: *Loch et al., ADNDT, **92** 813 (2006)*
 - Be: *Loch et al., ADNDT, **94** 257 (2008)*
- We just recently finished the GCR data for B.

Boron isonuclear sequence data sources

- **Dielectronic Recombination**

- ✓ B^{4+} : *Badnell et al., Astron. Astrophys. 447 389 (2006)*
- ✓ B^{3+} : *Bautista et al., Astron. Astrophys. 466 755 (2007)*
- ✓ B^{2+} : *Colgan et al., Astron. Astrophys. 417 1183 (2004)*
- ✓ B^+ : *Colgan et al., Astron. Astrophys. 412 597 (2003)*

- **Excitation**

- ✓ B^{4+} : RMPS – *Ballance et al., J. Phys. B 36 3707 (2003)*

[We have performed our own $n=7$ calculation that was used instead of this file]

- ✓ B^{3+} : RMPS – *Ballance (unpublished) – available at ADAS*
- ✓ B^{2+} : RMPS – *Griffin et al., J. Phys. B 33 1013 (2000)*
- ✓ B^+ : RMPS – *Badnell et al., J. Phys. B 36 1337 (2003)*
- ✓ B: RMPS – *Ballance et al., J. Phys. B 40 1131 (2007)*

- **Ionization**

- ✓ B^{4+} : RMPS+DW – *Griffin et al., J. Phys. B 38 L199 (2005)*
- ✓ B^{3+} : CCC + Expt – *Renwick et al., J. Phys. B 42 175203 (2009)*
- ✓ B^{2+} : RMPS/TDCC/CCC – *Badnell & Griffin., J. Phys. B 33 2955 (2000)*
- ✓ B^+ : TDCC, RMPS, DW + Expt – *Berregut et al., Phys. Rev. A 78 012704 (2008)*
- ✓ B: TDCC, RMPS – *Berregut et al., Phys. Rev. A 76 042704 (2007)*
- ✓ Excited states for B, B^+ and B^{2+} – *Lee et al., Phys. Rev. A 82 042721 (2010)*

The process of generating GCR data

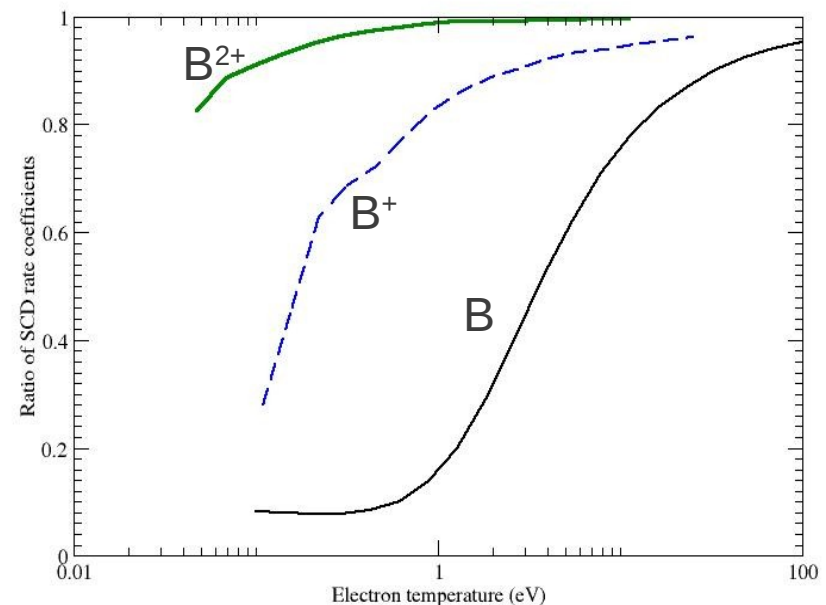
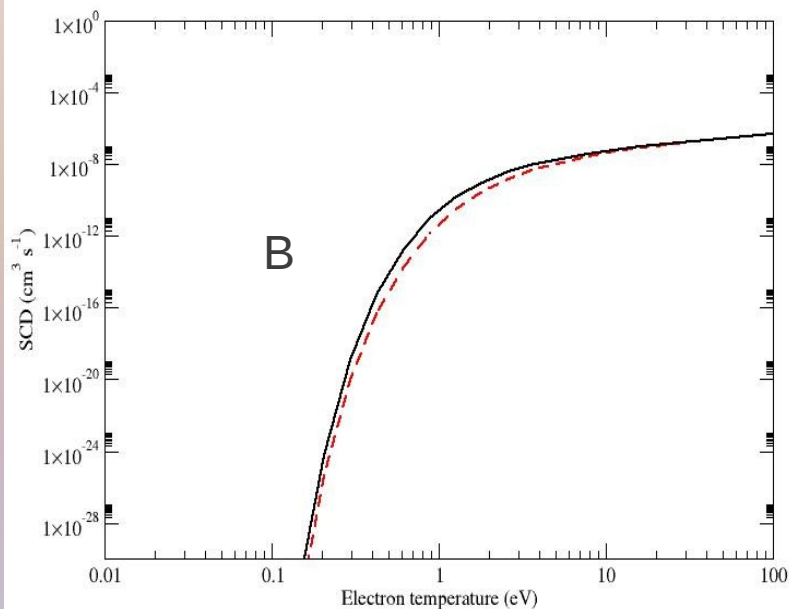
- Start with excitation datafile (R-matrix)
 - Supplement with non-dipole A-values
- Add RR+DR data
- Add ionization data
- Generate data for high n-shells (projection matrix)
- Process data through ADAS collisional-radiative modeling codes.

B GCR results

- Note that B has the following metastables, so has quite a few metastable cross coupling coefficients.
 - B $2s^2 2p$ (2P), $2s 2p^2$ (4P)
 - B⁺ $2s^2$ (2S), $2s 2p$ (3P)
 - B²⁺ $1s^2 2s$ (2S)
 - B³⁺ $1s$ (1S), $1s 2s$ (3S)
 - B⁴⁺ $1s$ (2S)
- The data is needs some final checks, but should be released shortly.

B GCR results : ionization

- The excited states contribute significantly to the effective ionization.
- Note that the relative size of the excited state contribution decreases as one goes to higher charge states.



The state of carbon GCR

Dielectronic Recombination

C^{5+} : *Badnell et al., Astron. Astrophys.* **447** 389 (2006)

C^{4+} : *Bautista et al., Astron. Astrophys.* **466** 755 (2007)

C^{3+} : *Colgan et al., Astron. Astrophys.* **417** 1183 (2004)

C^{2+} : *Colgan et al., Astron. Astrophys.* **412** 597 (2003)

C^+ : *Altun et al., Astron. Astrophys.* **420** 775 (2004)

Excitation

C^{5+} : RMPS – *Ballance et al., J. Phys. B* **36** 3707 (2003)

C^{4+} : RMPS – *Loch & Ballance, (unpublished)* – available at ADAS

C^{3+} : RMPS, TDCC + DW – *Griffin et al., J. Phys. B* **33**, 1013 (2000)

C^{2+} : RMPS – *Mitnik et al., J. Phys. B* **36** 717 (2003)

C^+ : Work in progress.

C: Yang et al. PRA 87 012704 (2013)

Ionization

C^{3+} :RMPS– *Badnell & Griffin., J. Phys. B* **33**, 2955 (2000)

C^{2+} : TDCC, CCC, RMPS, DW + Expt – *Loch et al., Phys. Rev. A* **71**, 012716 (2005)

C^+ : TDCC, RMPS, DW + Expt –*Ludlow et al., Phys. Rev. A* **78**, 052708 (2008)

C: TDCC, DW + Expt – *Pindzola et al., Phys. Rev. A* **62**, 042705 (2000).

Excited states of C^{3+} – Pindzola et al., JPCS 388, 062016 (2012)

Excited states of C^+ - Ballance et al. PRA 84, 062713(2011)

Excited states of C - Abdel-Naby et al. PRA, 76022708 (2013)

Conclusions

- The GCR data for B is complete.
 - It requires a few more checks and then can be put into the databases.
- The next element on the list is carbon. The electron-impact excitation of C^+ needs to be done, along with the excited state ionization of C^{2+} . Then the GCR data can be generated.